

Prašina

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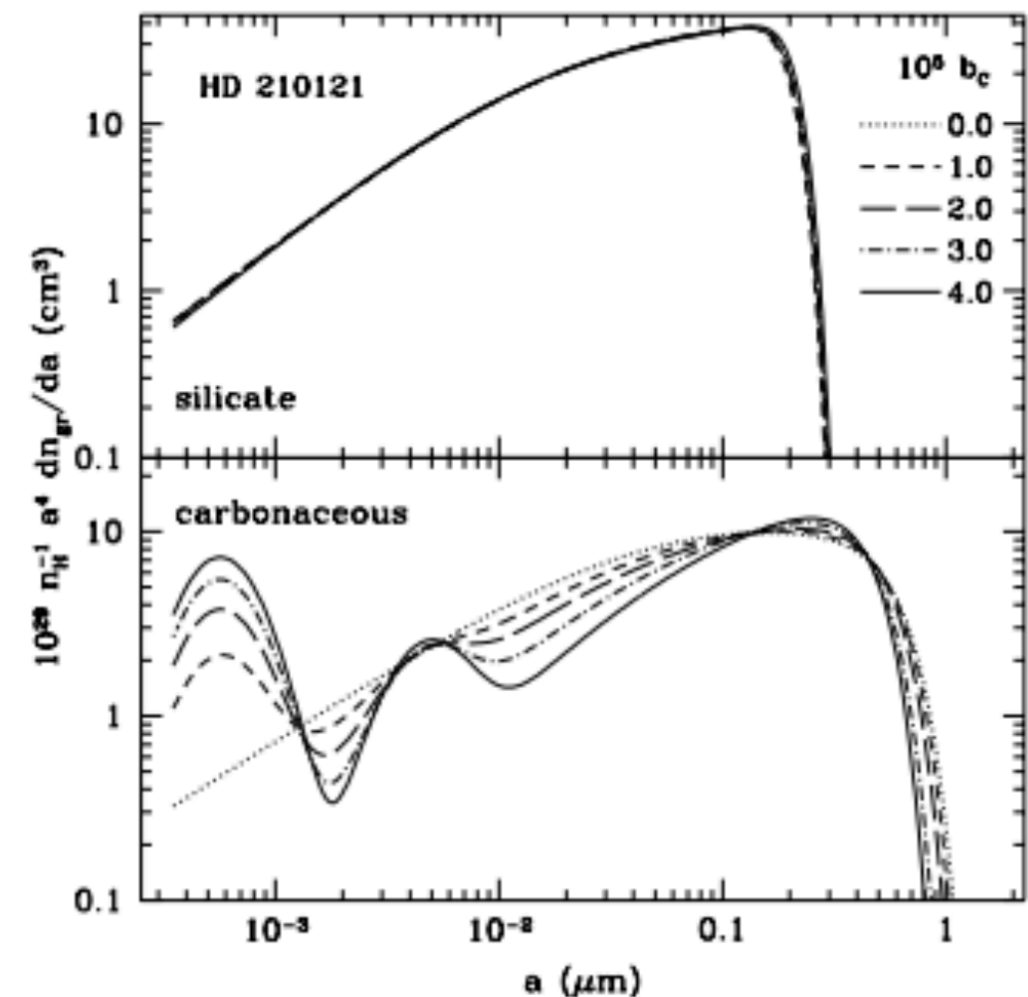
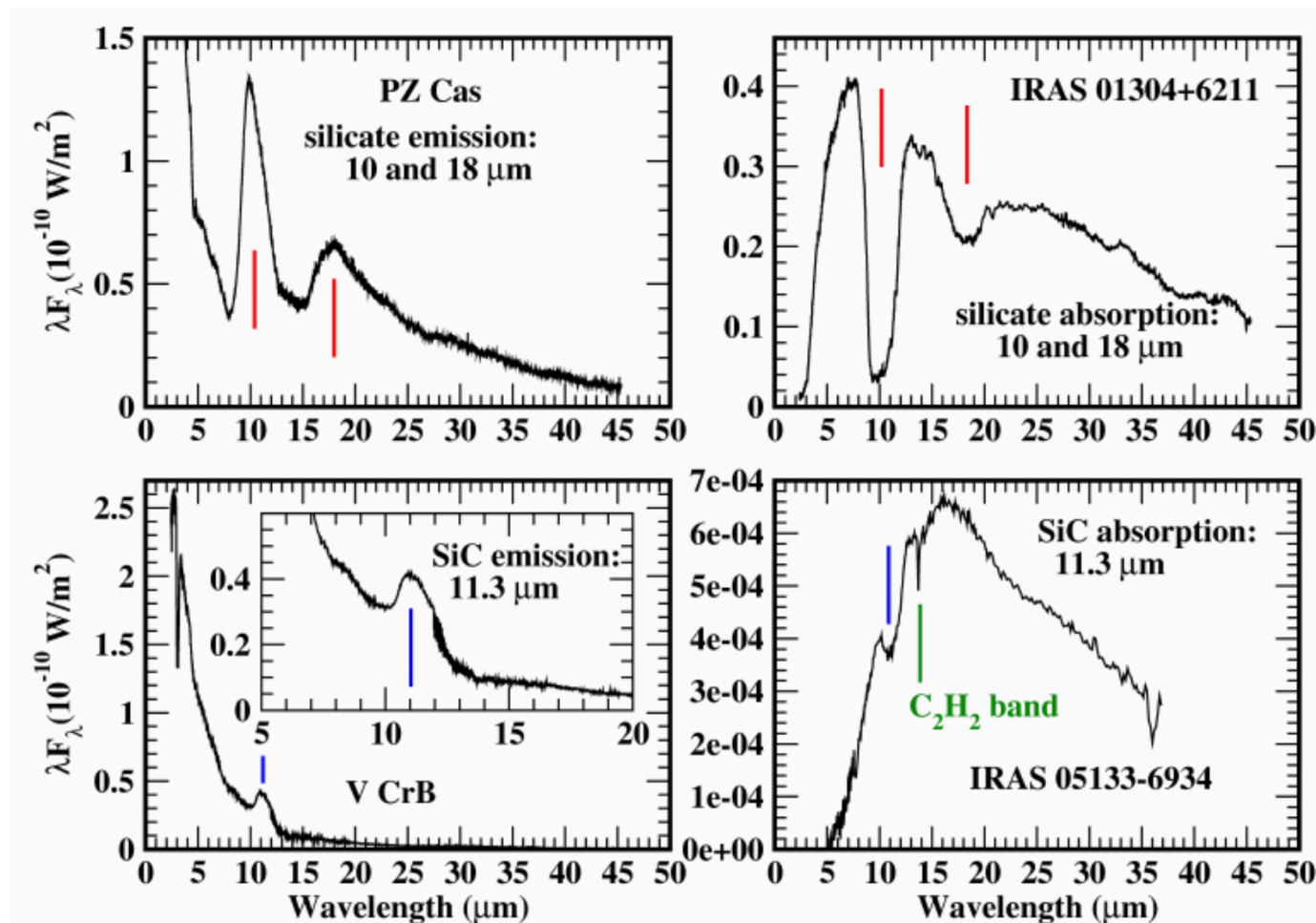
Osobine prašine

- Veličina utiče na:
 - apsorpciju, rasejanje svetlosti
 - emisiju od strane prašine
- Jezgro prašine utiče na
 - raspodelu po veličinama
 - izgled profila emisije/apSORpcije

zrno je optički debelo za:

$$a > \lambda$$

$$dn/da \sim a^{-3.5}$$



Ekstinkcija - zrno prašine

- Efektivni presek pojedinačnog zrna*: $\sigma = \pi a^2$
- Efektivni presek pri absorpciji: $\sigma_{\text{abs}} = \sigma Q_{\text{abs}}(\lambda)$
- Emitivnost^(?) (eng. *Emittance*): $\pi B_{\nu}(T) Q_{\nu}$
- Zrno u TRD se opisuje sa T_G
- Ukoliko se nalazi u polju izotropnog zračenja crnog tela na T (spec. slučaj - stvarnost: *razređena svetlost zvezda*)
 - iz svakog pravca zrno absorbira: $\sigma B_{\nu}(T) Q_{\text{abs}}(\lambda)$
 - u slučaju potpune ravnoteže ($T_G = T$) ukupno: $4\pi a^2 \pi Q_{\nu} B_{\nu}(T)$
- Važi ravnoteža: $Q_{\text{abs}}(\lambda) = Q_{\nu}$
- Više zrna?

Ekstinkcija zbog prašine

- Po jedinici zapremine:

$$j_\nu \rho = 1/(4\pi) \text{(number of grains per unit volume)} \text{(emission per grain)} \\ = \underline{N} \underline{4\pi a^2 \pi Q_\nu B_\nu(T)} / (4\pi) = N \pi a^2 Q_\nu B_\nu(T)$$

- Koeficijent po jedinici zapremine:

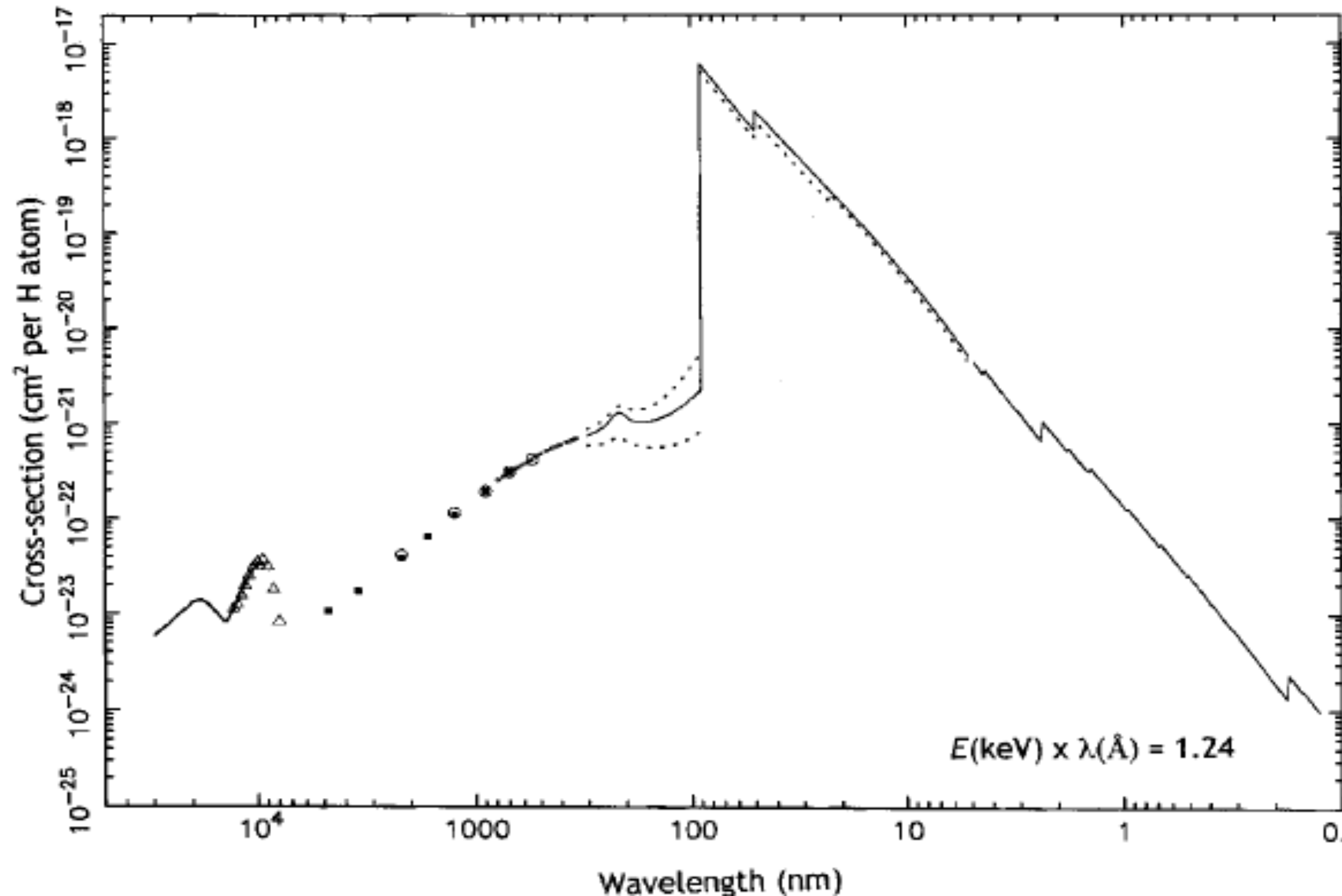
$$\kappa_\nu \rho = N \pi a^2 Q_\nu$$

- Maseni koeficijent absorpcije:

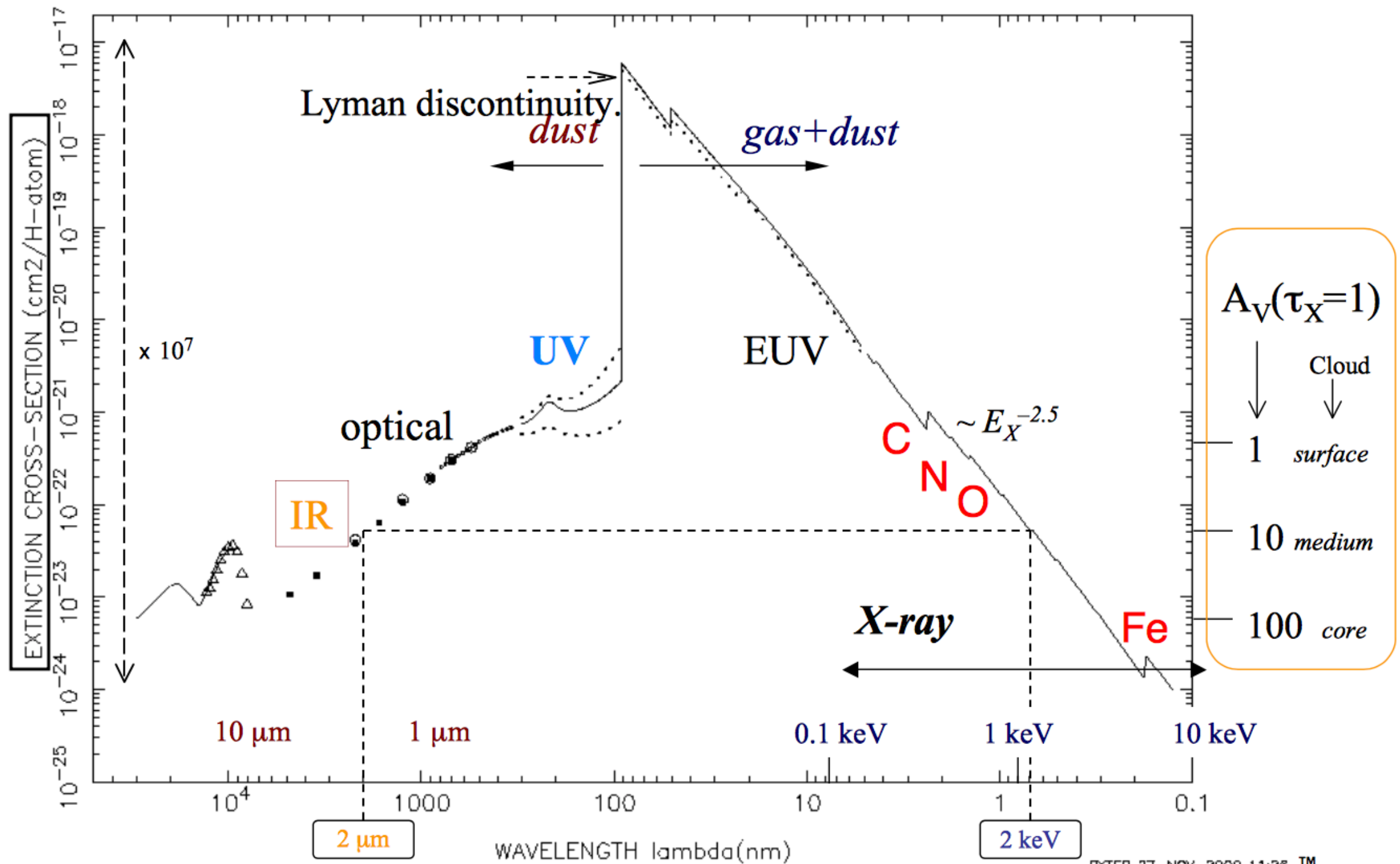
$$\kappa_\nu = \frac{N \pi a^2 Q_\nu}{N \frac{4}{3} \pi a^3 \rho_G} = \frac{3 Q_\nu}{4 a \rho_G}$$

- efikasnije za kraće talasne dužine ($\propto \lambda^{-1}$)
 - najefikasnije kada su sličnih dimenzija ($a \sim \lambda$)
- geometrijski poprečni presek \sim efektivni poprečni presek

Ekstinkcija (ukupno)



Ekstinkcija (detaljnije)



Ravnoteža?

- U MZM postoji ravnoteža između:
 - *sudara čestica, pritiska gasa, hemijskih reakcija...*
 - presudno: **apsorpcije i emisije**
- **Polje zračenja je opisano:**

$$I_\nu = WB_\nu(T_{is}) \quad I_\nu = c/(4\pi)u_\nu$$

- Ravnoteža je data sa:

$$4\pi\pi a^2 \int B_\nu(T_{is}) W Q_\nu d\nu = 4\pi a^2 \pi \int Q_\nu B_\nu(T_G) d\nu$$

$$W \int B_\nu(T_{is}) Q_\nu d\nu = \int Q_\nu B_\nu(T_G) d\nu$$

- Srednja vrednost:

$$\langle Q_\nu(T) \rangle = \frac{\int B_\nu(T) Q_\nu d\nu}{\int B_\nu(T) d\nu} = \frac{\pi}{\sigma_{SB} T^4} \int B_\nu(T) Q_\nu d\nu$$

$$W \frac{\sigma_{SB}}{\pi} T_{is}^4 \langle Q_\nu(T_{is}) \rangle = \frac{\sigma_{SB}}{\pi} T_G^4 \langle Q_\nu(T_G) \rangle$$

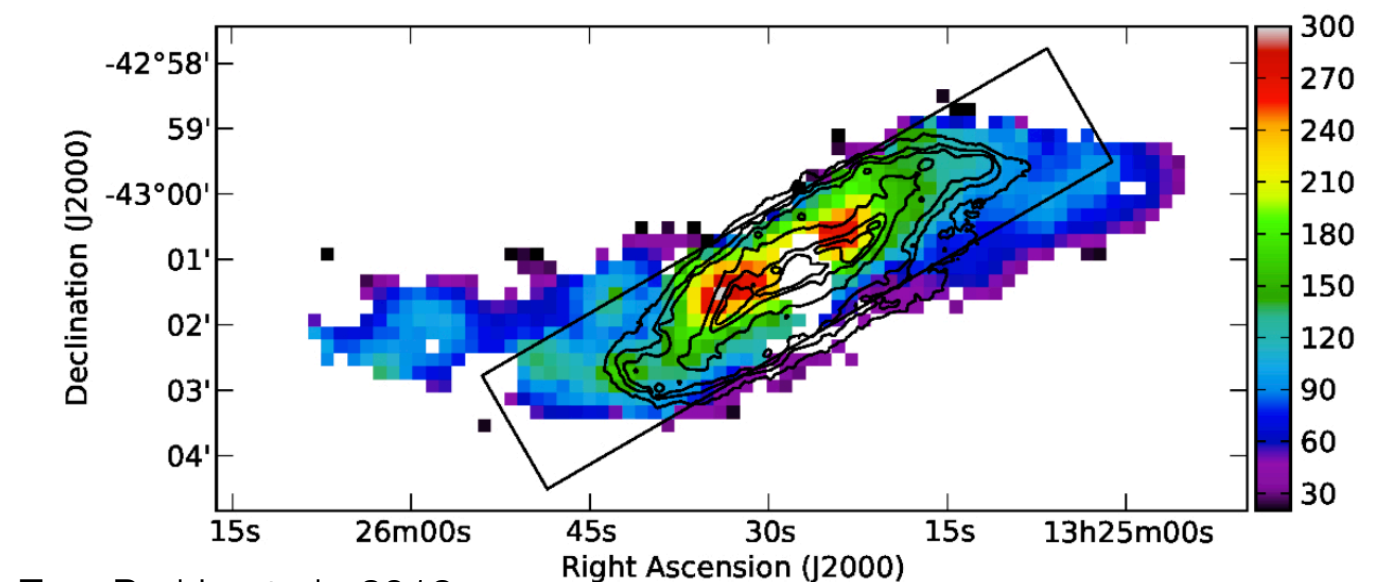
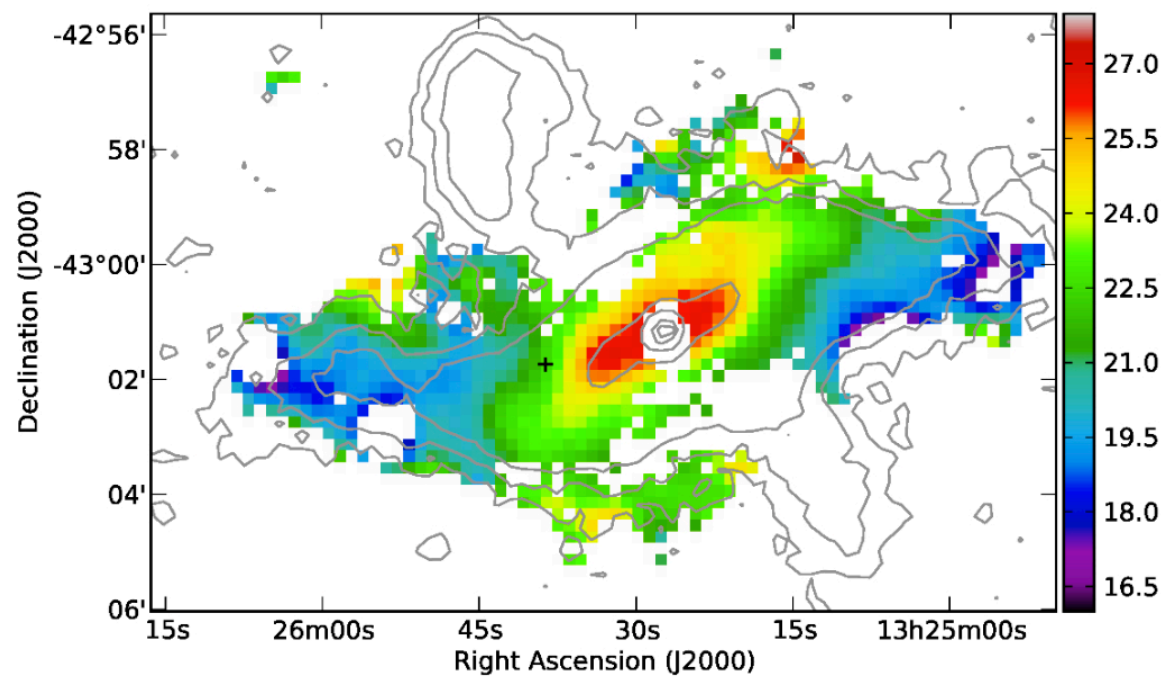
$$T_G = T_{is} W^{1/4} \left(\frac{\langle Q_\nu(T_{is}) \rangle}{\langle Q_\nu(T_G) \rangle} \right)^{1/4}$$

Veza između T_{is} i T_{G}

- U slučaju ravnoteže važi:

$$T_{\text{G}} \sim T_{\text{is}} W^{1/4} (\langle Q_{\text{UV}} \rangle / \langle Q_{\text{FIR}} \rangle)^{1/4}$$

- Za $W \sim 10^{-14}$ i odnos između $\langle Q_{\text{UV}} \rangle / \langle Q_{\text{FIR}} \rangle \sim 1$
 - T_{G} iznosi oko 3K
 - * ispravnost procene zavisi od:
 - * sastava zrna (menjaju Q)
 - * optičke debljine oblaka



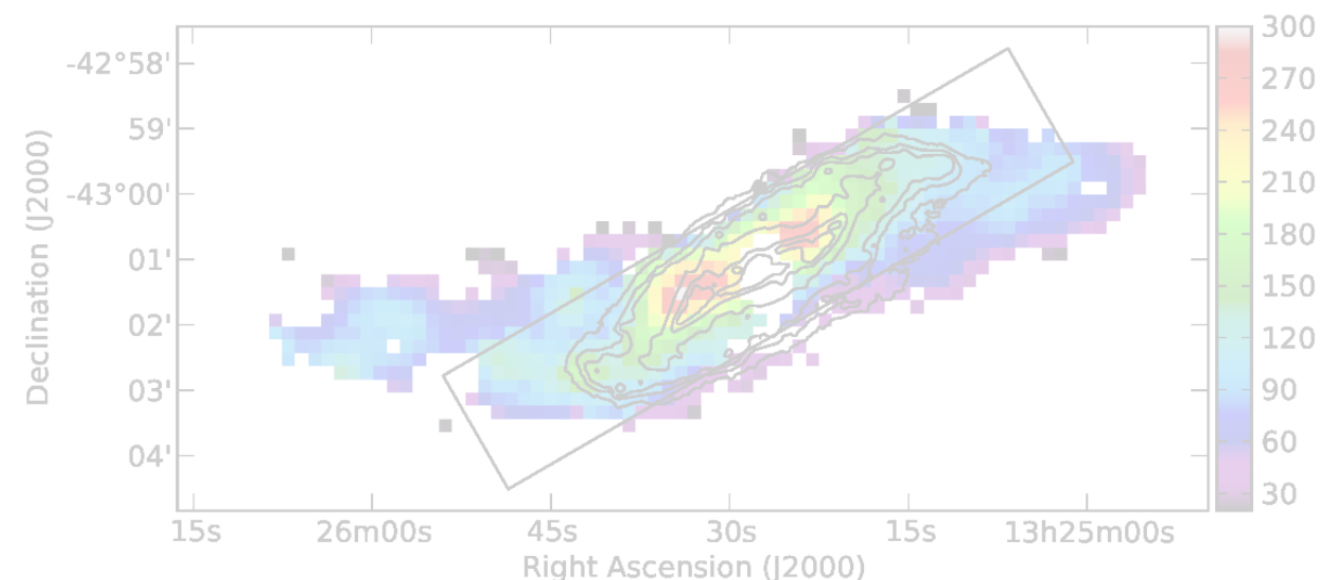
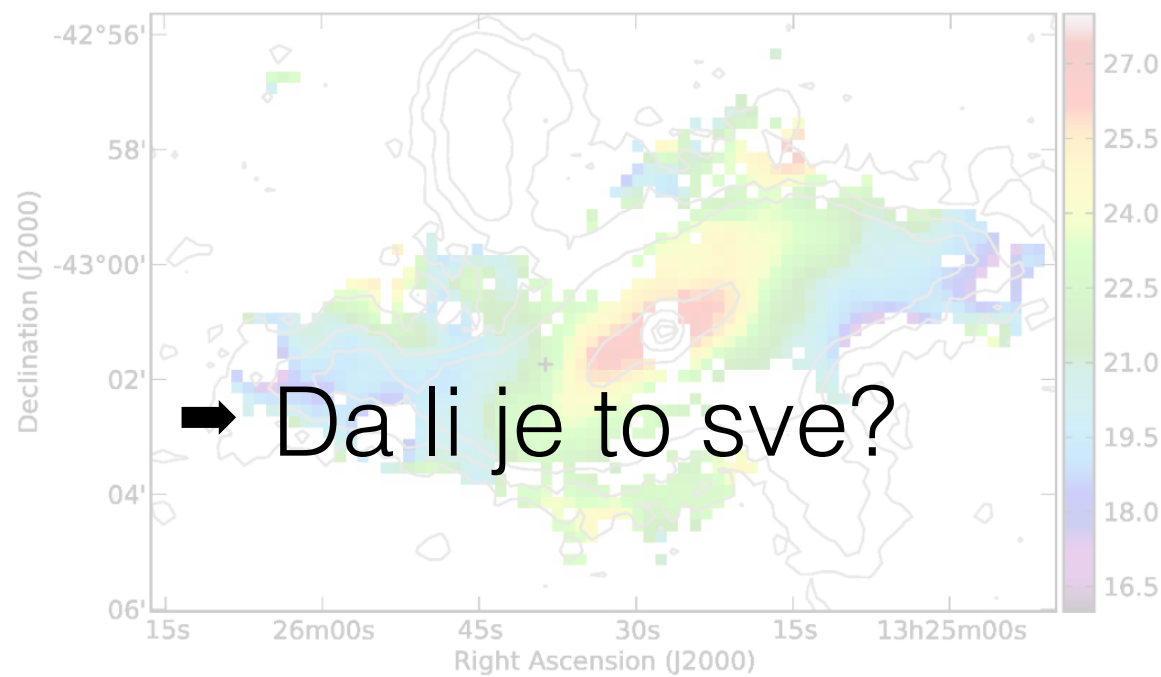
Tara Parkin et al., 2012

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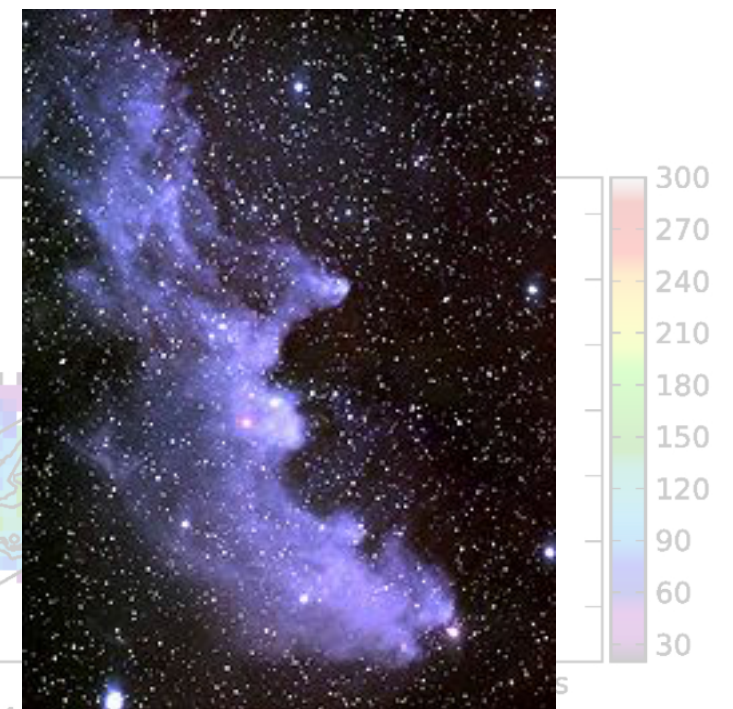
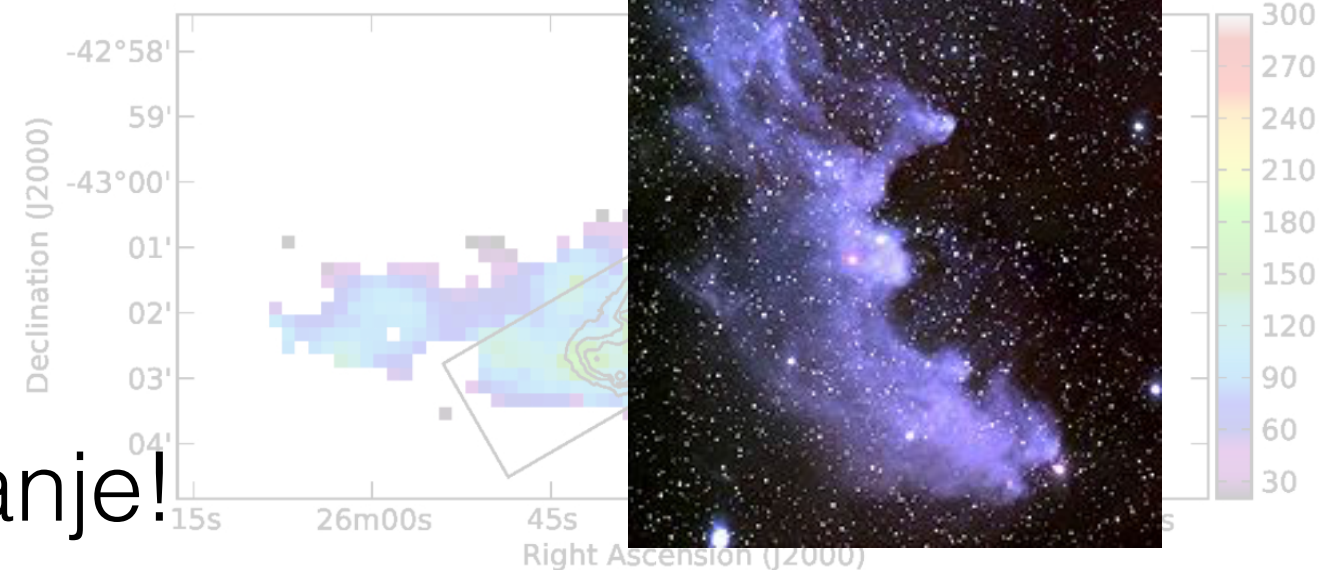
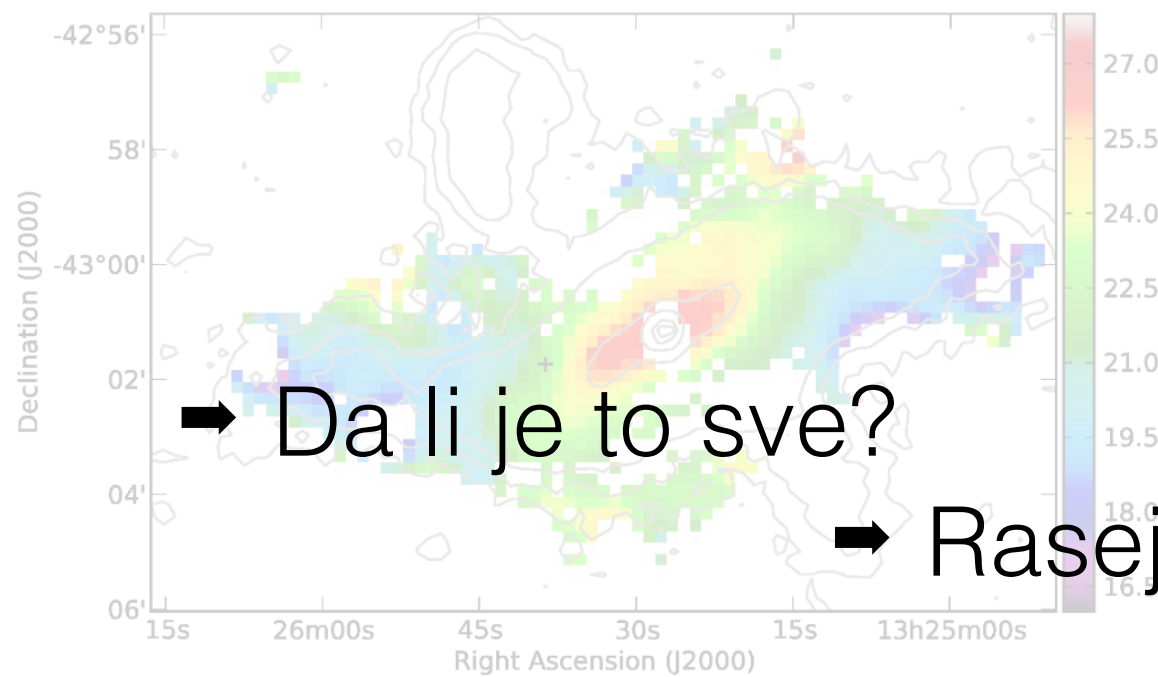


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Vez

s i T_G

- U slučaju

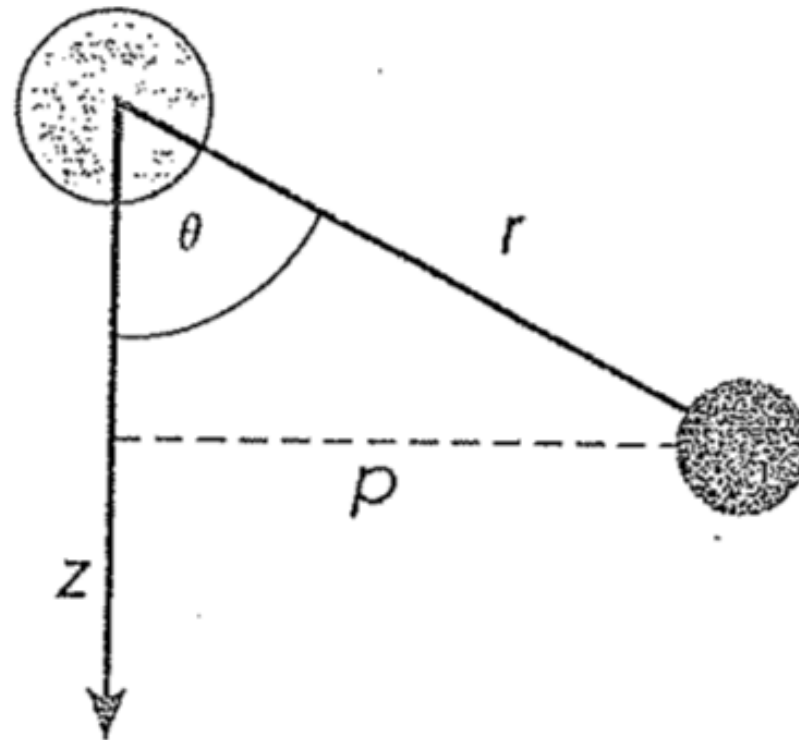
- Za $W \sim 10$

- T_G izm

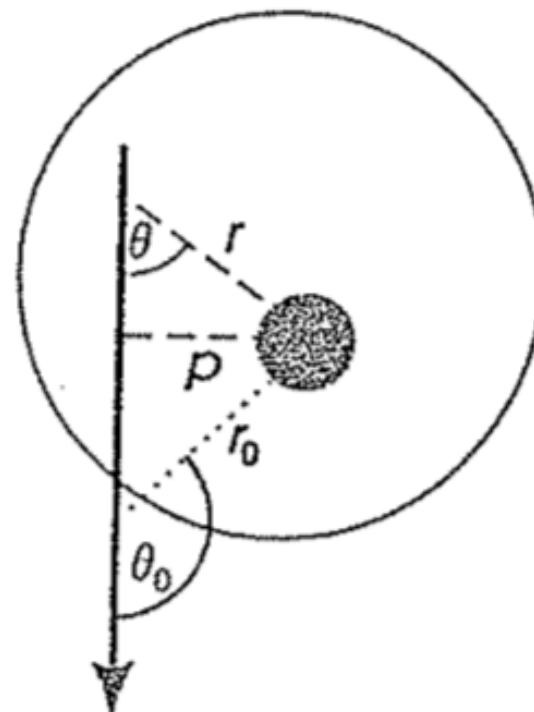
- * isprav

- * S

- * O



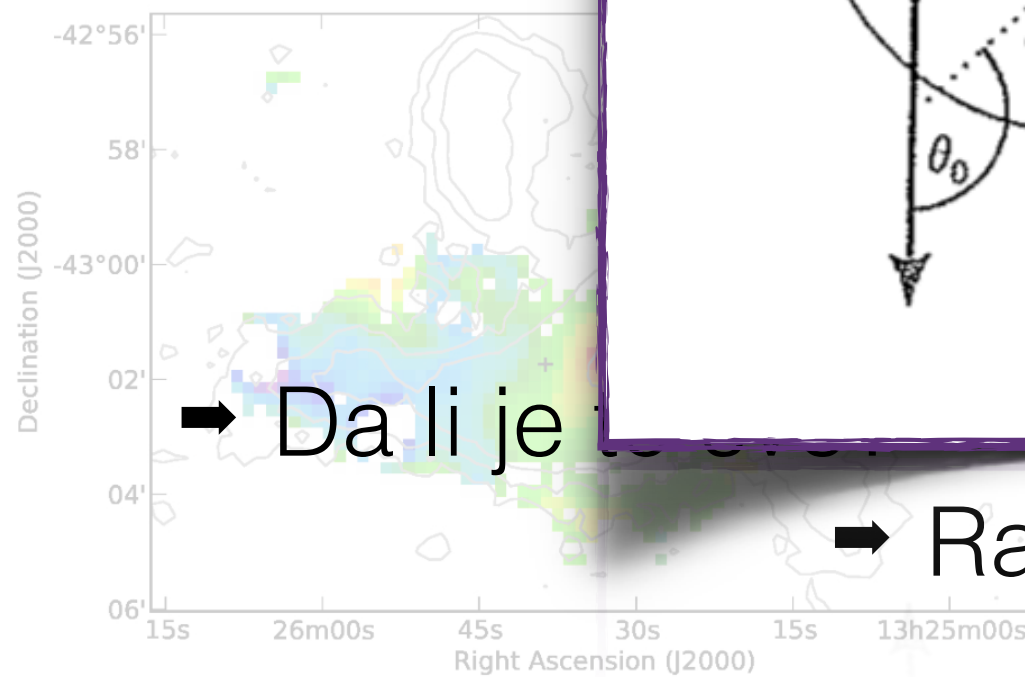
(a)



(b)

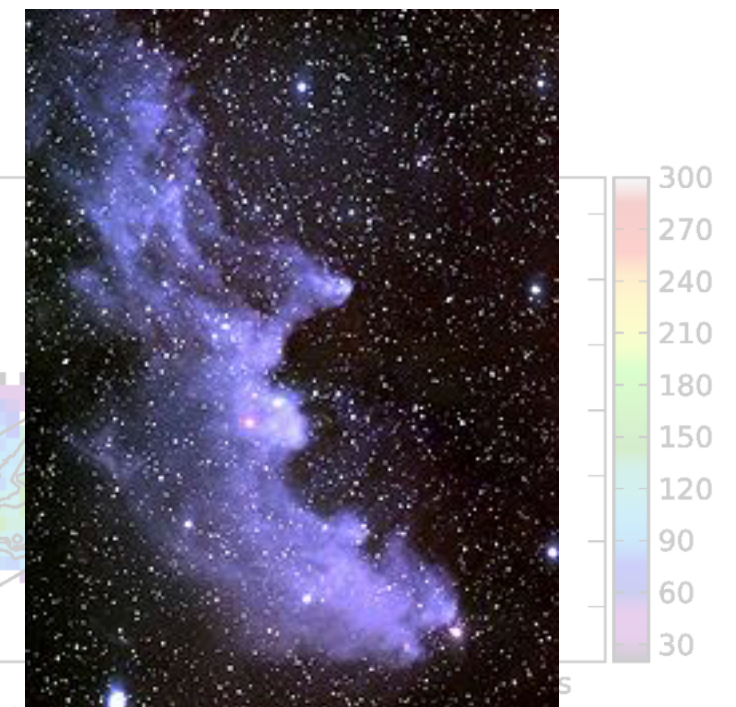
$$)^{1/4}$$

$$\langle \text{FIR} \rangle \sim 1$$

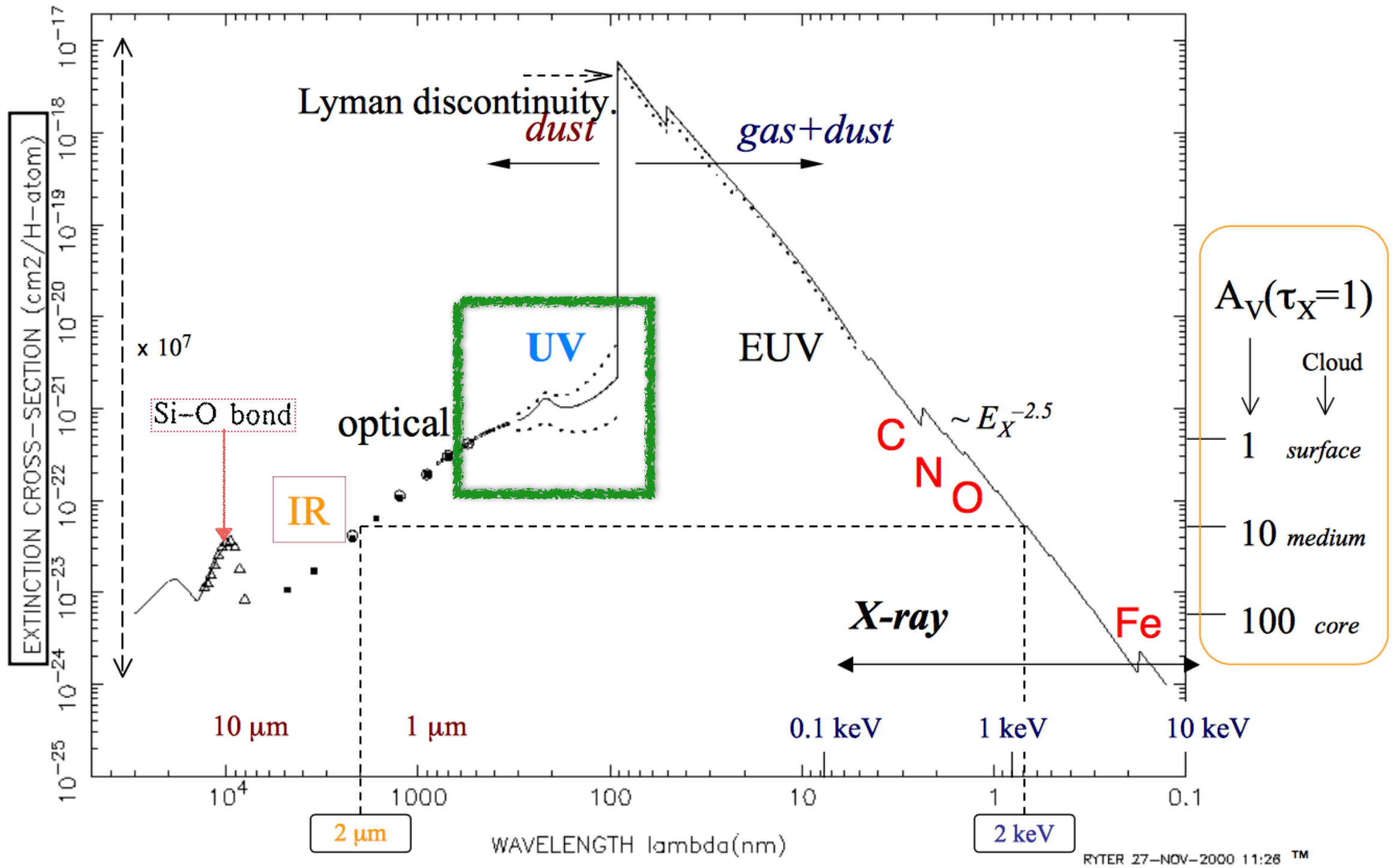


→ Da li je

→ Rasejanje!



Ekstinkcija (detaljnije)



T. Montmerle (2002)

I rasejanje?

- Ukupan koeficijent

$$K_{\text{ext}}\rho = K_{\text{abs}}\rho + K_{\text{sca}}\rho = N\sigma Q_{\text{ext}}(\lambda) = N\sigma[Q_{\text{sca}}(\lambda) + Q_{\text{abs}}(\lambda)]$$

$$\sigma_{\text{sca}} = \sigma Q_{\text{sca}}(\lambda)$$

- Albedo: $Q_{\text{sca}}/Q_{\text{ext}} = 1 - Q_{\text{abs}}/Q_{\text{ext}}$
- Koeficijenti $Q_{\text{sca}}(x)$ i $Q_{\text{abs}}(x)$ zavise od talasne dužine različito, ali po zakonu $Q = \text{constant } \nu^\beta$ gde $1 < \beta < 2$
- Bitan parametar je odnos između dimenzije zrna i tal. dužine $x = 2\pi a/\lambda$
- Mieva teorija - prašina opisana preko $m = \underbrace{n}_{\text{rasejanje}} - \underbrace{ik}_{\text{apsorpcija}}$

$$Q_{\text{abs}} \simeq 4x \Re \left[\frac{m^2 - 1}{m^2 + 1} \right] \quad Q_{\text{sca}} \simeq \frac{8}{3} x^4 \left| \frac{m^2 - 1}{m^2 + 1} \right|^2 \quad \text{slabo-apsorbujući}$$

$$m \sim n,$$

Koeficijent ekstinkcije

- Ukupan uticaj usled rasejanja i apsorpcije

- metod parova (dve zvezde iste klase)

$$\Delta m_\lambda = m_\lambda^{(1)} - m_\lambda^{(2)} = 5 \log \frac{r_1}{r_2} + \mathcal{A}_\lambda^{(1)} - \mathcal{A}_\lambda^{(2)}$$

$$\mathcal{A}_\lambda^{(i)} = 1.086 \tau_\lambda^{(i)} = 1.086 \int_0^{r_i} \kappa_\lambda \rho_{\text{ms}}(r'_i; l_i, b_i) dr'_i.$$

$$\Delta m(\lambda) = c_1 + c_2 \kappa(\lambda),$$

$$\Gamma(\lambda) = \frac{\Delta m(\lambda) - \Delta m(\lambda_2)}{\Delta m(\lambda_1) - \Delta m(\lambda_2)} = \frac{\kappa(\lambda) - \kappa(\lambda_2)}{\kappa(\lambda_1) - \kappa(\lambda_2)}$$

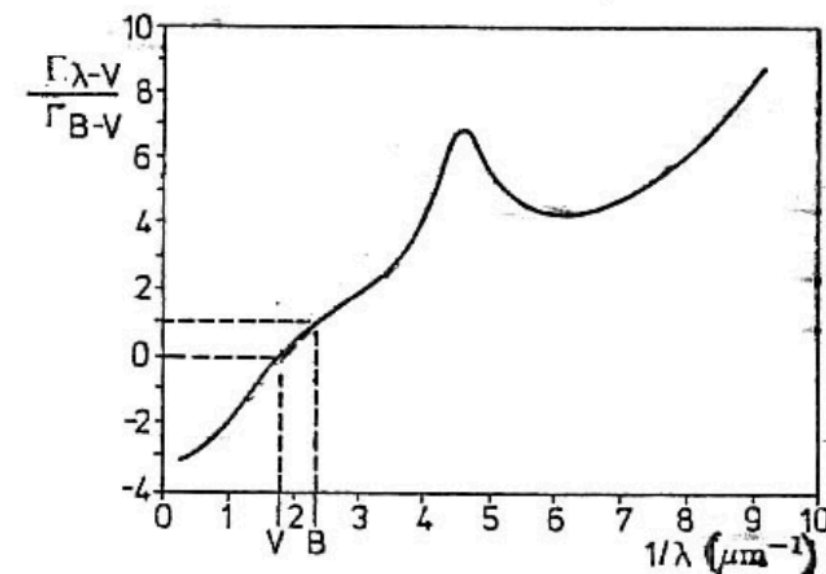
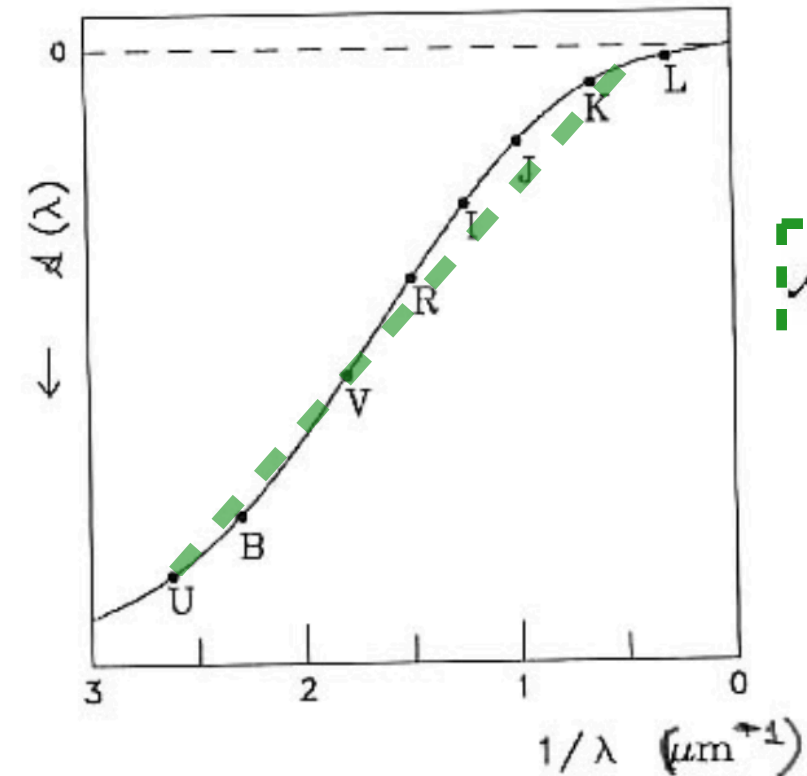
$$\mathcal{A}^{(2)} = 0 \text{ i } \mathcal{A}^{(1)} \equiv \mathcal{A}$$

$$E(\lambda_1, \lambda_2) = [m(\lambda_1) - m^0(\lambda_1)] - [m(\lambda_2) - m^0(\lambda_2)] = \mathcal{A}(\lambda_1) - \mathcal{A}(\lambda_2).$$

$$\Gamma(\lambda) = \frac{\mathcal{A}(\lambda) - \mathcal{A}(\lambda_2)}{\mathcal{A}(\lambda_1) - \mathcal{A}(\lambda_2)} = \frac{E(\lambda, \lambda_2)}{E(\lambda_1, \lambda_2)},$$

- c_1 i c_2 ne zavise od λ , ali zavise od r

- $\Gamma(\lambda)$ ne zavisi od r



Procena površinske gustine prašine

- Ukoliko N zrna veličine a izvrši ekstinkciju $Q_{\text{ext}}(\lambda)a^2N$,
- Ukoliko je linijska gustina za zrna ove veličine $n_{\text{col}}(a)da$
- Na osnovu posmatranja je poznato da u IC delu spektra važi $A(\lambda) \sim 1/\lambda^{1.84}$
- Integracijom $\Delta m = A(\lambda) = 2.5 \log_{10}[\exp(-\tau_D)] = 1.086 \tau_D = 1.086 \langle \pi a^2 Q_{\text{ext}}(a, \lambda) \rangle N_{\text{coll}}$
- Dobija se: $\int A(\lambda) d\lambda = 1.086 \int \pi a^2 n_{\text{col}}(a) \int Q_{\text{ext}}(a, \lambda) d\lambda da$
- Primenom Mieove teorije rasejanja imamo: $\int Q_{\text{ext}}(a, \lambda) d\lambda = 4\pi^2 a \frac{m^2 - 1}{m^2 + 2}$
- Što dalje daje: $\int A(\lambda) d\lambda = 1.086 * 4\pi^3 \frac{m^2 - 1}{m^2 + 2} \int a^3 n_{\text{col}}(a) da$

$$= 1.086 \times \frac{3\pi^2}{\rho_G} \frac{m^2 - 1}{m^2 + 2} \int m_G(a) n_{\text{col}}(a) da, \quad \text{where } m_G = \frac{4}{3} \pi a^3 \rho_G$$

$$= 1.086 \times \frac{3\pi^2}{\rho_G} \frac{m^2 - 1}{m^2 + 2} M_{\text{dust}}$$

+ posmatranja vodonika
-> Dust to Gas ratio

$$M_{\text{dust}} = \frac{A_V \int \frac{A(\lambda)}{A_V} d\lambda}{1.086 \times 3\pi^2} \rho_G \frac{m^2 + 2}{m^2 - 1}$$

Posmatranja prašnjavih oblaka

$$L_{\text{sca}}(\text{UV}) \quad L_{\text{abs}}(\text{FIR})$$

- Odnos UV i FIR posmatranja može se iskoristiti kao mera albeda $w = L_{\text{sca}} / (L_{\text{sca}} + L_{\text{abs}})$

- Na pojedinačnm zrnu se raseje:

$$Q_{\text{sca}} \pi a^2 L_v / (4\pi r^2) \exp[-\tau_{\text{ext}}(r)] \quad L_v \sim \pi B_v(T^*) 4\pi R^2$$

- Za izotropno* rasejanje:

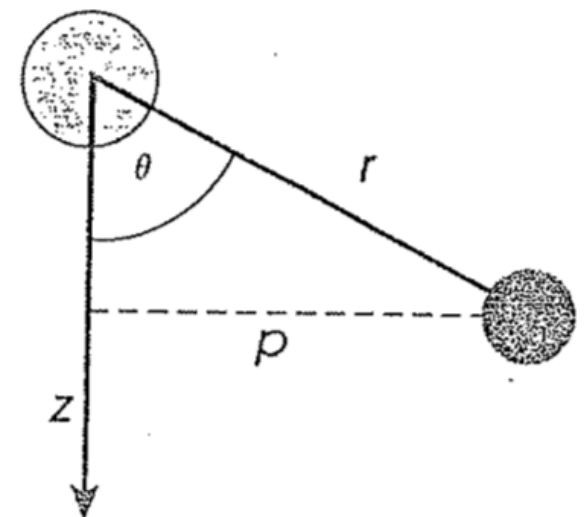
$$I_{\text{sca}} d\Omega = \frac{L_v}{4\pi r^2} e^{-\tau_{\text{ext}}} Q_{\text{sca}} \pi a^2 N f(\theta) dV d\Omega$$

- Na rastojanju p posmatramo:

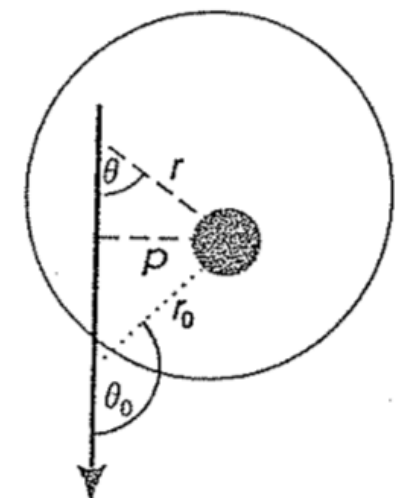
$$I(p) = \int Q_{\text{sca}} \pi a^2 \frac{L_v}{4\pi r^2} N e^{-\tau_{\text{ext}}(r)} e^{-\tau_{\text{ext}}(z,p)} f(\theta) dz$$

$$= \omega \int \frac{L_v}{4\pi r^2} e^{-\tau_{\text{ext}}(r)} e^{-\tau_{\text{ext}}(z,p)} f(\theta) d\tau_{\text{ext}}(z)$$

$$d\tau_{\text{ext}}(z) = 1/\omega d\tau_{\text{sca}} = 1/\omega N Q_{\text{sca}} \pi a^2 dz$$



(a)



(b)

Posmatranja prašnjavih oblaka

$$L_{\text{sca}}(\text{UV}) \quad L_{\text{abs}}(\text{FIR})$$

- Odnos UV i FIR posmatranja može se iskoristiti kao mera albeda $L_{\text{sca}}/(L_{\text{sca}} + L_{\text{abs}})$

- Naprednija metoda posmatranja

$$Q_{\text{sca}} \tau$$

- Za

Kako je oblak daleko manjih dimenzija od rastojanja do njega ($r \sim \text{const}$)

$$I(p) = \omega \frac{L_v}{16\pi^2 r^2} e^{-\tau_{\text{ext}}(r)} (1 - e^{-\tau_{\text{ext}}})$$

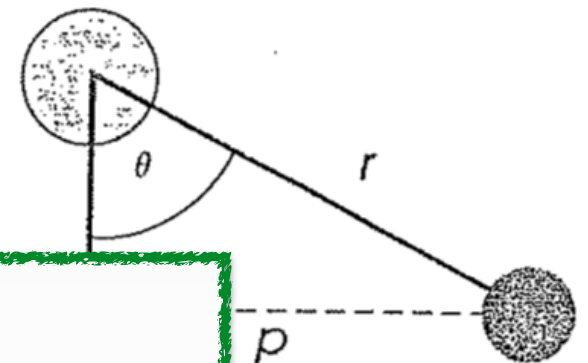
- N

$$F_v(*) = L_v/(4\pi d^2) \exp[-\tau_{\text{ext}}(\text{SO})]$$

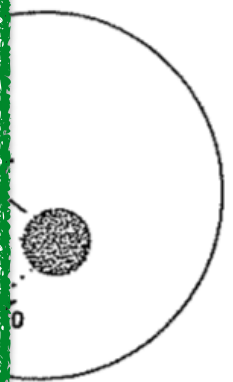
$$I(p) = \int Q_{\text{sca}} \pi a^2 \frac{1}{4\pi r^2} I_0 e^{-\tau_{\text{ext}}(r)} e^{-\tau_{\text{ext}}(z, p)} f(\theta) dz$$

$$= \omega \int \frac{L_v}{4\pi r^2} e^{-\tau_{\text{ext}}(r)} e^{-\tau_{\text{ext}}(z, p)} f(\theta) d\tau_{\text{ext}}(z)$$

$$d\tau_{\text{ext}}(z) = 1/\omega d\tau_{\text{sca}} = 1/\omega N Q_{\text{sca}} \pi a^2 dz$$



(a)

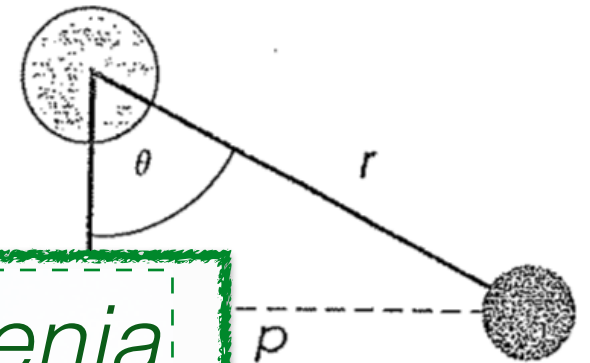


(b)

Posmatranja prašnjavih oblaka

$$L_{\text{sca}}(\text{UV}) \quad L_{\text{abs}}(\text{FIR})$$

- Odnos UV i FIR posmatranja može se iskoristiti kao mera albeda $L_{\text{sca}}/(L_{\text{sca}} + L_{\text{abs}})$



- Napomena: $+ ima još rešenja$

- Kako je oblak daleko manjih dimenzija od rastojanja do njega ($r \sim const$)

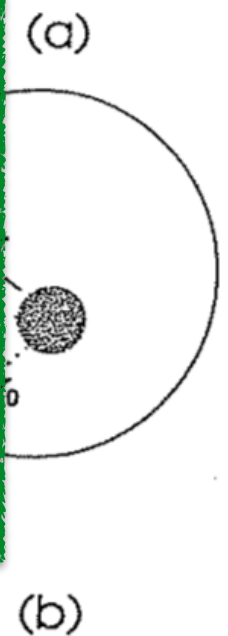
$$I(p) = \omega \frac{L_v}{16\pi^2 r^2} e^{-\tau_{\text{ext}}(r)} (1 - e^{-\tau_{\text{ext}}})$$

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$$I(p) = \int Q_{\text{sca}} \pi a^2 \frac{1}{4\pi r^2} I_v e^{-\tau_{\text{ext}}(r)} e^{-\tau_{\text{ext}}(z, p)} f(\theta) dz$$

$$= \omega \int \frac{L_v}{4\pi r^2} e^{-\tau_{\text{ext}}(r)} e^{-\tau_{\text{ext}}(z, p)} f(\theta) d\tau_{\text{ext}}(z)$$

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Ostala rešenja

- U slučaju sfernog oblaka sa zvezdom u sredini:
 - ka centru oblaka $\tau_{\text{ext}}(r) = \kappa \rho p / \sin \theta$
 - prema obodu oblaka $\tau_{\text{ext}}(z, p) = \kappa \rho [\sqrt{r_0^2 - p^2} + p \cot \theta]$

$$\begin{aligned}
 I(p) &= \frac{\omega \kappa_{\text{ext}} \rho L_v}{4\pi} \int_{\theta_0}^{\pi - \theta_0} \frac{1}{r^2} f(\theta) \exp \left\{ -\kappa_{\text{ext}} \rho \left[\frac{p}{\sin \theta} + \sqrt{r_0^2 - p^2} + p \cot \theta \right] \right\} d(p \cot \theta) \\
 &= \frac{\omega \tau_0 L_v}{4\pi r_0} \int_{\theta_0}^{\pi - \theta_0} f(\theta) \frac{\sin^2 \theta}{p^2} \exp \left\{ -\kappa_{\text{ext}} \rho \left[\frac{p}{\sin \theta} + \sqrt{r_0^2 - p^2} + p \cot \theta \right] \right\} \frac{-p}{\sin^2 \theta} d\theta
 \end{aligned}$$

... dalje, ukoliko uzmemo $dz = d(p \cot \theta)$ i $\tau_0 = \kappa_{\text{ext}} \rho r_0$,
 + uz pps optički retke sredine važi:

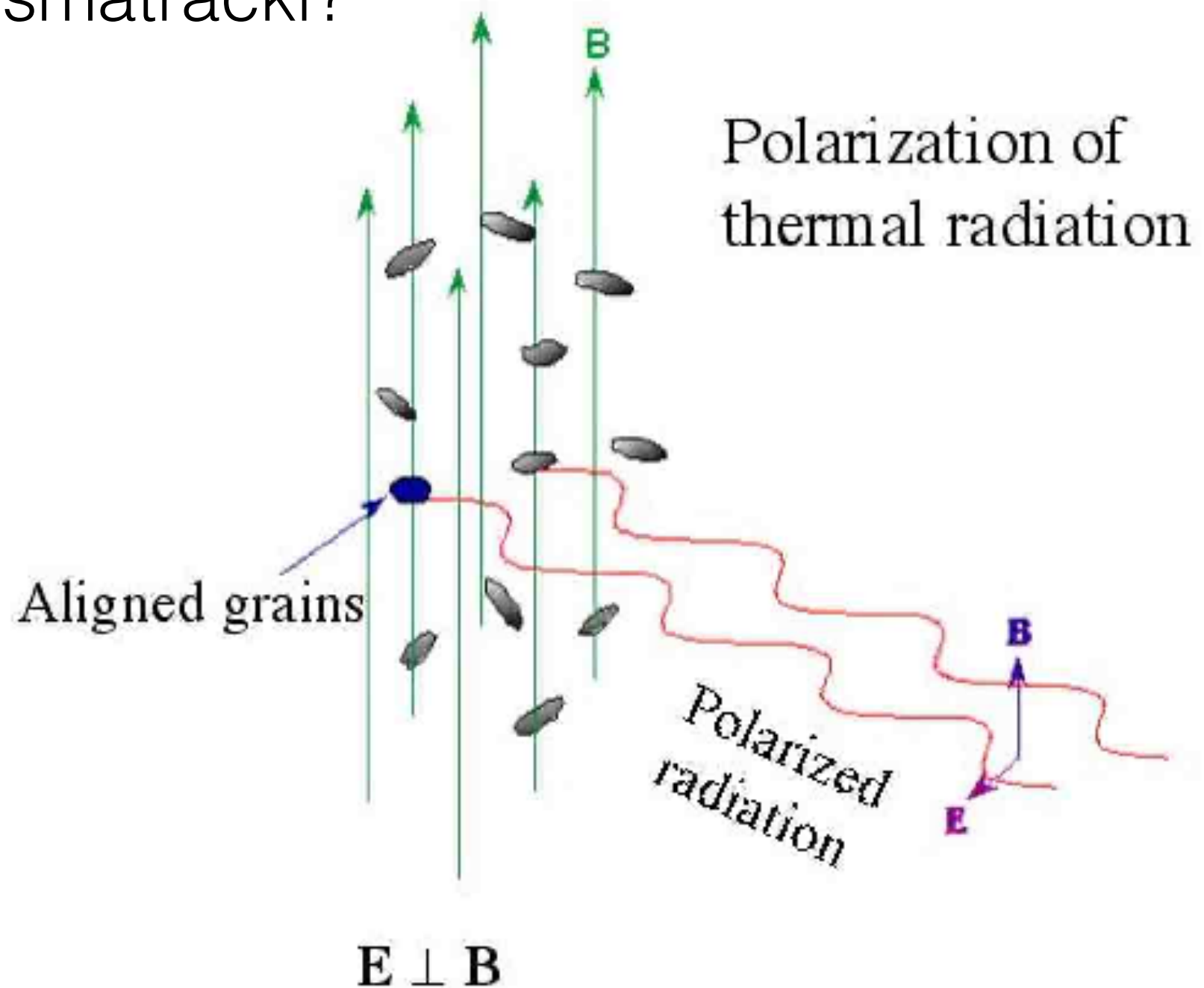
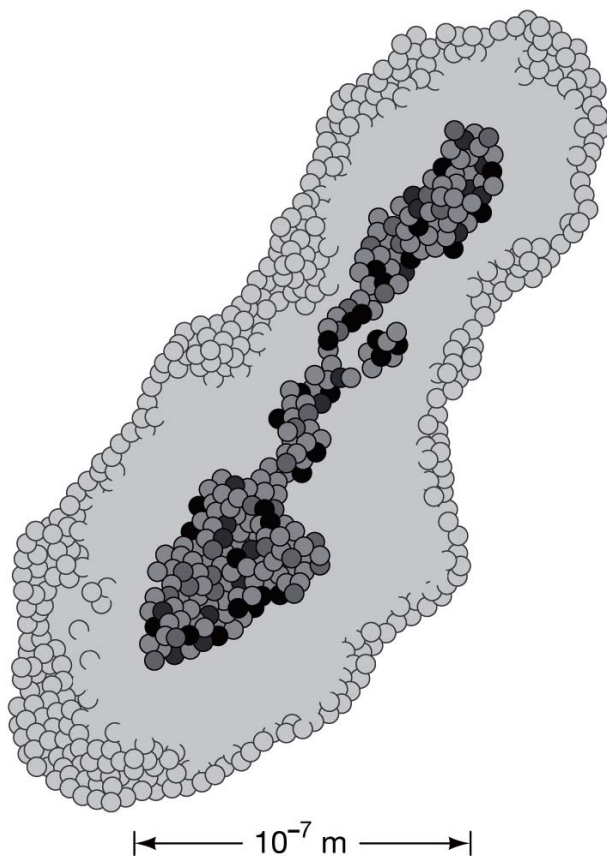
$$\begin{aligned}
 \frac{d[pI(p)]}{dp} &= \frac{\omega \tau_0 L_v}{4\pi r_0^2 \cos \theta_0} \frac{d}{d\theta_0} \int_{\theta_0}^{\pi - \theta_0} f(\theta) d\theta, \quad \text{since } dp = r_0 \cos \theta \\
 &= \frac{\omega \tau_0 L_v}{4\pi r_0^2} \frac{f(\theta_0) + f(\pi - \theta_0)}{\cos \theta_0}
 \end{aligned}$$

$$I(p) = \frac{\omega L_v \tau_0}{4\pi r_0 p} \int_{\theta_0}^{\pi - \theta_0} f(\theta) d\theta$$

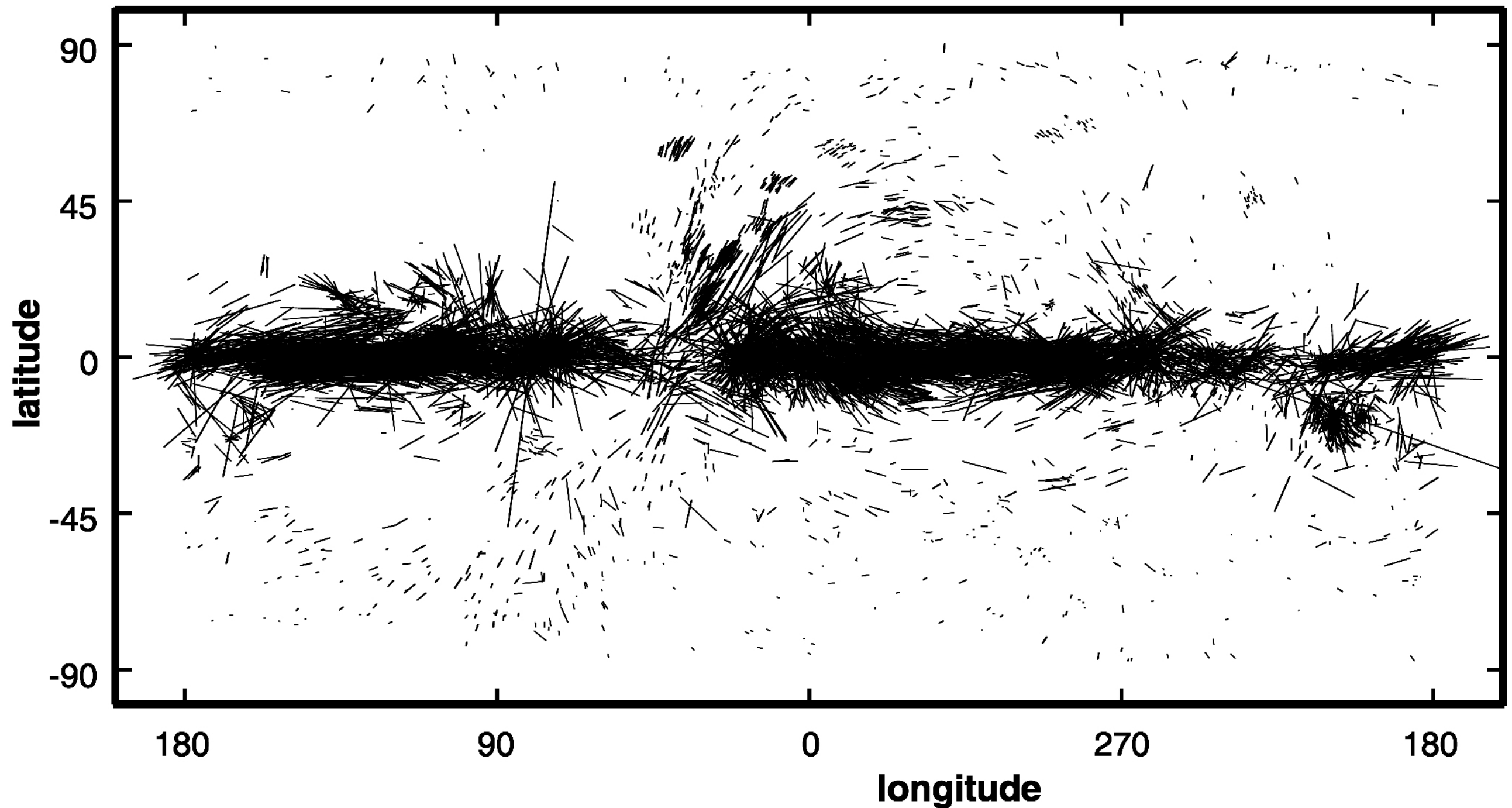
aproximativno

Polarizacija reflektovane svetlosti

- Sferna simetrija zrna \neq dobra pretpostavka
- Kako ovo izgleda posmatrački?

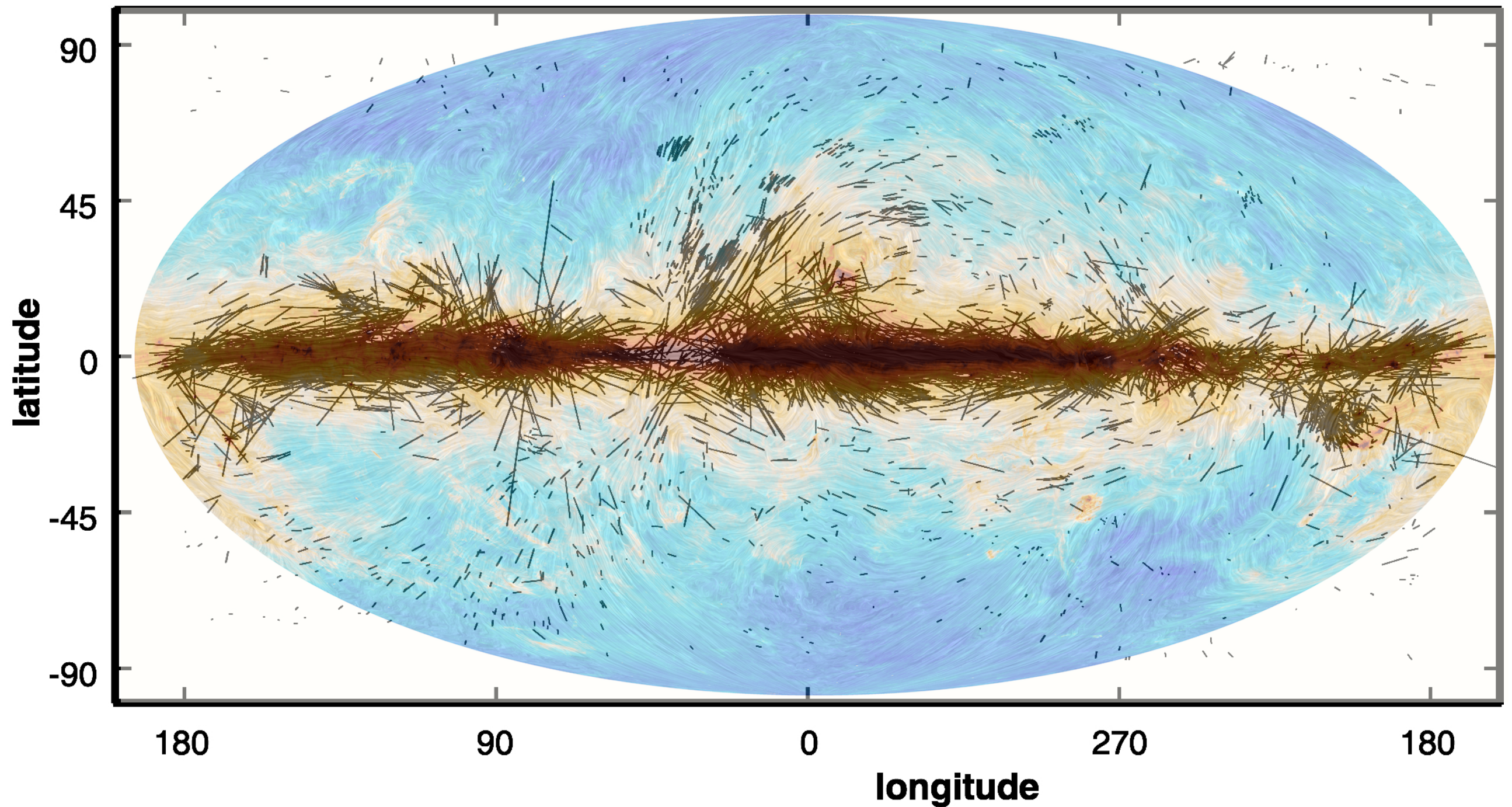


Polarizacija reflektovane svetlosti



(Berdyugin et al. 2014)

Polarizacija reflektovane svetlosti



(Planck collaboration 2016)

(Berdyugin et al. 2014)

Hvala na pažnji!

